

Particle Physics in the Early Universe (FS 17)

EXERCISE SERIES 7

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Exercise 1 In this exercise we will show that in a Grand Unified Theory with a single coupling, that has the SM group as a subgroup, there is a condition between the coupling constants stemming from pure group theory normalization. Let us define the set of all SM fermions of a single family as

$$S = \{u_L^i, d_L^i, u_R^i, d_R^i, e_L, e_R, \nu_L\} \quad (1)$$

where $i = R, G, B$, the color quantum number of the quarks

- (i) QCD generators are normalised usually as $\tilde{T}^i = \frac{1}{2}g_s\lambda^i$ where λ^i are the Gell-Mann matrices. Consider the generator

$$\tilde{T}^3 = \frac{1}{2}g_s\lambda^3 \quad , \quad \lambda_3 = \text{diag}(+1, -1, 0) \quad (2)$$

Show that the sum over all particles in S of the squares of the eigenvalues of \tilde{T}^3 is equal to

$$4 \cdot \left(\left(\frac{g_s}{2} \right)^2 + \left(-\frac{g_s}{2} \right)^2 \right) = 2g_s^2 \quad (3)$$

- (ii) $SU(2)_L$ generators are also normalised as $T^i = \frac{1}{2}g\sigma^i$, where σ^i are the Pauli matrices. Show that the sum over all particles in S of the squares of the eigenvalues of T^3 is equal to

$$2 \cdot 3 \cdot \left(\left(\frac{g}{2} \right)^2 + \left(\frac{g}{2} \right)^2 + \left(-\frac{g}{2} \right)^2 \right) = 2g^2 \quad (4)$$

- (iii) The $U(1)$ eigenvalue, Y , of each particle are normalised so that $Y = T_3 - Q$, where $T_3 = \pm 1, 0$ depending on whether the particle is in the upper or lower part of an $SU(2)$ doublet, or whether it is an $SU(2)$ singlet, respectively. Show that the sum over all particles in S of the square of the eigenvalue of the $U(1)$ generator, $g'Y$ is equal to

$$\dots = \frac{10}{3}g'^2 \quad (5)$$

Requiring that the three sums are equal at and above the unification scale provides the condition

$$g_s(M) = g(M) = \sqrt{\frac{5}{3}}g'(M) \quad (6)$$

Exercise 2 The coupling of a gauge theory depends on the typical scale at which the interaction takes place. The dependence is determined by the Renormalization Group Equation (RGE), that, in general, has the form

$$\mu \frac{dg}{d\mu} = -\frac{g^3}{4\pi^2}\beta_1 + \dots \quad (7)$$

where we ignore higher order terms that appear because of renormalization beyond the one-loop level. Integrate this equation from a low scale μ to the GUT scale M , to get

$$\frac{1}{g^2(\mu)} = \frac{1}{g^2(M)} - \frac{\beta_1}{2\pi^2} \log\left(\frac{M}{\mu}\right) \quad (8)$$

There is one such equation for each SM coupling, g_s, g, g' , with different constants $\beta_{1s}, \beta_{1w}, \beta_{1Y}$. We will use

$$\beta_{1s} = \frac{11}{4} - \frac{n_g}{3} \quad \beta_{1w} = \frac{11}{6} - \frac{n_g}{3} \quad \beta_{1Y} = -\frac{5n_g}{9} \quad (9)$$

where n_g is the number of generations, now known to be equal to 3.

Assuming unification at scale M , i.e. that $g_s(M) = g(M) = \sqrt{\frac{5}{3}}g'(M)$, show that, setting $\mu = m_z$ we can get

$$\log \frac{M}{m_z} = \frac{4\pi}{11e^2(m_z)} \left(1 - \frac{8e^2(m_z)}{3g_s^2(m_z)}\right) \quad (10)$$

Does this results depend on the number of generations, n_g ? Finally, assume the RG equation is known one order higher, that is

$$\mu \frac{dg}{d\mu} = -\frac{g^3}{4\pi^2}\beta_1 - \frac{g^5}{(4\pi^2)^2}\beta_2 + \dots \quad (11)$$

Integrate it to find $g(\mu)$.

Reminder: $e = g \sin \theta_W = g' \cos \theta_W$