## Particle Physics in the Early Universe (FS 17)

## Exercise Series 2

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**Exercise 1** For any given metric  $g_{\mu\nu}$ , the Christoffel symbols,  $\Gamma^{\mu}_{\nu\rho}$  satisfy

$$\Gamma^{\mu}_{\nu\rho} = \frac{1}{2} g^{\mu\sigma} \left( \vartheta_{\nu} g_{\sigma\rho} + \vartheta_{\rho} g_{\sigma\nu} - \vartheta_{\sigma} g_{\nu\rho} \right) \tag{1}$$

Show that, for a general Robertson-Walker type metric, that is, any metric of an isotropic homogeneous space, defined by

$$d\tau^{2} = g_{\mu\nu}dx^{\mu}dx^{\nu} = dt^{2} - a(t)^{2}\gamma_{ij}dx^{i}dx^{j}$$
(2)

where

$$\gamma_{ij} = h(r)dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$
(3)

and

$$h(r) = \frac{1}{1 - \tilde{k}r^2} \quad , \quad \tilde{k} \equiv \frac{k}{L^2} \tag{4}$$

(i) the Christoffel symbols with two or more indices equal to zero vanish:

$$\Gamma_{0i}^0 = \Gamma_{00}^0 = 0 \tag{5}$$

(ii) the Christoffel symbols with one index equal to zero are given by

$$\Gamma_{ij}^0 = a\dot{a}\gamma_{ij} \tag{6}$$

and

$$\Gamma^i_{0j} = \frac{\dot{a}}{a} \delta^i_j \tag{7}$$

Reminder:  $\dot{a} \equiv \frac{da(t)}{dt}$ 

(iii) the Christoffel symbols with all indices different than zero are independent of a(t). They are identical with the Christoffel symbols of the metric  $\gamma_{ij}$ :

$$\Gamma_{jk}^{i} = \frac{1}{2} \gamma^{im} \left( \vartheta_{j} \gamma_{mk} + \vartheta_{k} \gamma_{mj} - \vartheta_{m} \gamma_{jk} \right) \tag{8}$$

Exercise 2 Let's look at the Christoffel symbols with space indices. Show by direct computation and using step 3 of exercise 1 that the only non-vanishing Christoffel symbols are

$$\Gamma_{rr}^r = \tilde{k} \, r h(r) \tag{9}$$

$$\Gamma^{i}_{rj} = \frac{1}{r} \delta^{i}_{j} \quad , \quad i, j \in \{\theta, \phi\}$$
 (10)

$$\Gamma_{ij}^{r} = -\frac{r}{h(r)}\tilde{\gamma}_{ij} \quad , \quad i, j \in \{\theta, \phi\} \quad \tilde{\gamma}_{ij} = diag(1, \sin^{2}\theta)$$
 (11)

$$\Gamma^{\theta}_{\phi\phi} = -\sin\theta\cos\theta \tag{12}$$

$$\Gamma^{\phi}_{\theta\phi} = \frac{\cos\theta}{\sin\theta} \tag{13}$$

Exercise 3 The Ricci tensor is defined as

$$R_{\nu\rho} = \vartheta_{\mu} \Gamma^{\mu}_{\nu\rho} - \vartheta_{\nu} \Gamma^{\mu}_{\mu\rho} + \Gamma^{\mu}_{\mu\sigma} \Gamma^{\sigma}_{\nu\rho} - \Gamma^{\mu}_{\nu\sigma} \Gamma^{\sigma}_{\mu\rho} \tag{14}$$

(i) Using the results of the previous exercise, show that

$$R_{00} = -3\frac{\ddot{a}}{a} \tag{15}$$

and

$$R_{ij} = \left(a\ddot{a} + 2\dot{a}^2 + 2\tilde{k}\right)\gamma_{ij} \tag{16}$$

(ii) Compute the curvature scalar  $R \equiv g^{\mu\nu}R_{\mu\nu}$  and show that the left-hand side of Einstein's equations for the 0-0 component is

$$R_{00} - \frac{1}{2}Rg_{00} = 3\left(\frac{\dot{a}^2}{a^2} + \frac{\tilde{k}}{a^2}\right) \tag{17}$$