

Exercise 1. Phase transitions in the electroweak theory

We will consider the electroweak theory a $SU(2)_L \times U(1)_Y$ broken by a Higgs scalar. The Lagrangian density of this model has the form

$$\mathcal{L} \supset \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}} \quad (1)$$

where the gauge part consists of the usual field strength tensors:

$$\mathcal{L}_{\text{gauge}} = -\frac{1}{4}W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4}B_{\mu\nu}B^{\mu\nu}, \quad a = 1, 2, 3 \quad (2)$$

with the $SU(2)_L$ $W_{\mu\nu}^a \equiv \partial_\mu W_\nu^a - \partial_\nu W_\mu^a - g\varepsilon^{abc}W_\mu^b W_\nu^c$ and $U(1)_Y$ $B_{\mu\nu} \equiv \partial_\mu B_\nu - \partial_\nu B_\mu$. The Higgs is now an $SU(2)_L$ doublet,

$$H = \begin{pmatrix} H^+ \\ H^0 \end{pmatrix} \quad (3)$$

with Lagrangian density

$$\mathcal{L}_{\text{Higgs}} = (D_\mu H)^\dagger (D^\mu H) - m^2 H^\dagger H - \lambda (H^\dagger H)^2 \quad (4)$$

and covariant derivative

$$D_\mu H = \left(\partial_\mu + ig \frac{\sigma^a}{2} W_\mu^a + ig' \frac{1}{2} B_\mu \right) H. \quad (5)$$

The fermions acquire their mass through their Yukawa couplings to the Higgs field:

$$\mathcal{L}_{\text{Yukawa}} = - \sum_f y_f F_L^T i \sigma^2 H f_R + \text{h.c.} \quad (6)$$

where y_f denotes the Yukawa coupling, $F_L \equiv \begin{pmatrix} u_f \\ d_f \end{pmatrix}_L$ denotes the left-handed doublet, and f_R the right-handed singlet.

In the electroweak theory the symmetry is spontaneously broken, by the Higgs field acquiring a VEV

$$\langle H \rangle = \begin{pmatrix} 0 \\ \varphi_c / \sqrt{2} \end{pmatrix} \quad (7)$$

(a) Compute the masses of the fermions as a function of φ_c .

(b) Defining the photon, Z and W^\pm by

$$A_\mu \equiv \cos \theta_W B_\mu + \sin \theta_W W_\mu^3, \quad Z_\mu \equiv -\sin \theta_W B_\mu + \cos \theta_W W_\mu^3, \quad W_\mu^\pm \equiv W_\mu^1 \pm i W_\mu^2 \quad (8)$$

compute their masses as a function of φ_c .

When fermions are involved the result from the last exercise sheet needs to be adapted as follows:

$$\bar{V}_1^{T=0}(\varphi_c) = B' \varphi_c^4 \left[\ln \frac{\varphi_c^2}{v'^2} - \frac{25}{6} \right], \quad (9)$$

$$\bar{V}_1^{T \neq 0}(\varphi_c) = -\frac{\pi^2 T^4}{90} (n_S + n_V + \frac{7}{8} n_F) + \frac{T^2}{24} [t_S + 3t_V + 2t_F] - \frac{C'T}{3}, \quad (T \gg 0) \quad (10)$$

where the the unprimed parameters are the same as previously.

(c) With this result compute the 1-loop effective potential $\bar{V}(\varphi)$ and show that

$$m^2(T) = m^2 + \left(\frac{\lambda}{2} + \frac{e^2(1 + 2 \cos^2 \theta_W)}{4 \sin^2 2\theta_W} + \sum_f \frac{y_f^2}{12} \right) T^2 \quad (11)$$

(d) Retaining only the gauge boson contribution to the φ_c^3 -term in $\bar{V}(\varphi)$ show that

$$C' \approx \frac{3e^3(1 + 2 \cos^3 \theta_W)}{4\pi \sin^3 2\theta_W} \quad (12)$$

Dropping the λ^2 contributions to the radiative corrections one can prove that

$$B' \approx \frac{3}{4} \left(\frac{e^2}{4\pi} \right)^2 \frac{1 + 2 \cos^4 \theta_W}{\sin^4 2\theta_W} - \frac{1}{64\pi^2} \sum_f y_f^2. \quad (13)$$

(e) Find the sign of B' with the values of the parameters in the standard model ($e^2/4\pi \approx 1/137, \dots$). Write the numerical values and comment on which terms are essentially important and which ones are not.